

Calculations of Muon Radiation

for the Meson area

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Abstract

The muon radiation pattern produced by protons striking the M-center target has been studied at some length—using the C. E. R. N. computer program "HALO" (Ref. 1). The purpose of these studies has been to gain insight into the muon radiation pattern likely when targetting 1000 GeV protons. In the studies to-date positive pion production at the target is followed by decay and transport of the resulting muons through an approximation to the Meson area tunnels and shielding. The resulting radiation density is very sensitive to the presence of magnetic field in the first M2 bending magnets. Levels calculated due to $\frac{1}{1000}$ $\frac{1}{10$

Introduction

A prior study of muon fluxes in the Meson experimental area was made by T. Ferbel; it is described in Volume 2 of the 1976 Fermilab Summer Study. In comparing 1000 Gev production to 400 Gev production, a prediction of an increase by a factor of 100 was made. $\,$ For 3 $_{ exttt{X}}$ 10**12, 400 Gev protons on a 20 cm Be target the muon flux density in a 4 mr cone about the direction of the primary proton beam was calculated to be 1.5 x 10**5 $\,\mu$ /sq. meter (\sim .2 mr/hr for a 1976 accelerator cycle time of 8 sec). This 1976 report has served as a starting point for the present study. We have proceeded from it to include the effect of magnetic field downstream of the target; we have used a Monte-Carlo approach rather than a closed-form calculation; we have tried another production model; and we have more accurately approximated the shielding material and voids. Left for later study are: (1) "direct" muon production, (2) effects due to primary-energy protons striking material downstream of the target, (3) variations of the magnetic field arrangement downstream of the target, (4). refinements to the representation of the passive shielding in the area, and (5) investigation of other possible production models. Other topics for study would be the muons resulting from interactions in the M-West target and the muons that would result from targetting for the M1 High Intensity Pion beam.

Review of 1976 Ferbel Summer Study

In the course of the Ferbel Summer Study fits to π^+ , π^- , κ^+ , κ^- production data were made to cover the x (Feynman x) region of interest and these were used to calculate the number of decay muons above 200 GeV in energy. An estimate for multiple scattering was made in order to calculate the area density of these muons in the region of the M3 experiments. A 20 cm Be target was used—followed by a 5 meter decay path.

The production cross sections used were

$$\frac{d\sigma}{dx}(\pi^{+}) = 2.5 \frac{d\sigma}{dx}(\pi^{-})$$

$$\frac{d\sigma}{dx}(K^{+}) = 2 \frac{d\sigma}{dx}(K^{-})$$

and

$$\frac{d\sigma}{dx}(\pi) = 250 e^{-1/x} mb$$

$$\frac{d\sigma}{dx}(\kappa) = 22 e^{-10x} mb$$

both for x>0.2. These were the results of fits to data T. Ferbel had available and were good fits between 0.2 < x < 0.6 .

We have tried to reproduce the summer study results. For muons from pion parents and for 400 GeV incident protons we succeed in repeating the number of muons calculated to within 35% of Ferbel's number. For Kaon parents we come closer and get to within 6%. When we do the calculations for an incident energy of 1000 GeV and take the ratio to our numbers for 400 GeV incident energy, we get a ratio of 52. This number compares well to the number 60 quoted in the summer study report. This ratio is for muons with energy above 200 GeV and does not include the increase due to the decrease in the multiple scattering angle. Our calculation is further discussed in Appendix 1.

The summer study report quotes some measured values near the M3 beam for muon fluxes at 400 GeV and 300 GeV primary proton energies. These and the calculated numbers agreed fairly well (a factor of two). It is pointed out in the report that 4-10 feet of steel lowered measured rates by an order of magnitude.

Remarks on Production Models

The Stefanski-White production model (Ref. 2) has often been employed in Neutrino Department shielding studies and has been incorporated into a local version (Ref. 3) of the CERN program "HALO". It can be expressed for the production by 400 GeV protons as

$$\frac{d\sigma}{d\rho}$$
 (STEF-WHITE, π^{\dagger}) = 0.32 e mb

where

(see Appendix 2). The Ferbel 1976 S.S. production formula is

The x dependence of the two models is quite different. Using the Stefanski-White expression the ratio of $\mu + \mu$ production by 1000 GeV protons to $\mu + \mu$ production by 400 GeV protons for E($\mu + \mu$) > 200 GeV, is 17. This is much less than the equivalent number of 43 from the Ferbel model.

The comparison of these two, differing results leads to concern regarding an appropriate production model. Exploration of the question of what is the proper production model has been left out of the purview of this study. Between the Ferbel model and the Stefanski-White model it would be more conservative, when studying muon shielding, to use the Ferbel model—since it gives a higher prediction for 1000 Gev; however, this study has not been consistent in this usage.

Use of "HALO" to calculate muon fluxes

The CERN program "HALO" has been utilized to study muon fluxes in the Meson experimental area. This has been done with a succession of better and better approximations to the Meson area magnetic fields and passive steel and dirt shielding (see Fig. 1), for the case of the M-Center target that feeds the M1/M2/M3/M4 beamlines. The Stefanski-White production model has been used for the studies discussed in this section. Both 400 GeV and 1000 GeV were used as incident proton energies. The program was modified to take into account a vertical production angle of 0.6 milli-radian for the M3 beam — for the 400 GeV case. "HALO" has also been used to cross-check against the Ferbel results and these cross-checks are described in Appendix 3.

The Meson Area was first modelled for "HALO" without any magnetic fields present. The rates calculated near the M3 beam axis were high when compared with fluxes measured by the E533 group (see discussion in Appendix 4). Inclusion of magnetic fields in the calculation necessitated the generation of "HALO"-compatible magnetic field maps for the magnet of interest. We chose to do this on the Fermilab CDC Cyber 175 computers (Ref. 4). The magnets inserted were the first bend magnets in the M2 beam. These serve also as the first magnetic sweepers in the M3 line. The bend consists of 30 feet of B 2-type main ring bending magnets. The magnetic field is 18.5 kilogauss (25 mr bend for a 200 Gev particle). The z location starts at 236 feet. More details are given in Appendix 5.

With the M2/M3 magnets in place the resulting muon radiation pattern at z=1400 feet is shown in Figure 2. In this figure an entry of "1" represents a bin with a positive muon flux of 8.7 x 10**4 / per square meter. These results correspond to 10**12 400 Gev protons interacting in a "thin" target. The flux represented by a "1" entry in a bin in this figure corresponds to a dosage rate of 0.09 mr/hr (Ref. 5). Indicated on Fig. 2 is an area with the boundaries x -420 to -300 inches, y -20 to -4 inches. This area is centered in y on the intersection of the proton beam axis with the x,y plane shown (intersection at -10"). The average number of entries per bin in this area is 2.7; the corresponding dosage rate is 0.25 mr/hr (12 sec. cycle time).

For 1000 GeV protons (at 0 mr production angle) the muon radiation pattern at z=1400' is shown in Figure 3. In this case "1" represents a bin with flux 1.4 x 10**6 p^+ m⁻² — or 0.3 mr/hr (accelerator cycle is taken to be 60 sec). As before the normalization is 10**12 interacting protons. The area indicated on Fig. 3 has boundaries x -280 to -180 inches and y -8 to +8 inches. The average number of entries per hin in this area is 4.7; the corresponding dosage rate is 1.4 mr/hr. This represents an increase of a factor of 6 in going from 400 GeV to 1000 GeV. It includes a factor of 1/5 due to the longer accelerator cycle. A production model with a steeper x dependence (such as that of the 1976 Ferbel S.S.) would make this ratio higher than 6.

Muon Measurements

Muon measurements have been made to compare with the "HALO" calculations. These are described in detail in Appendix 6 and the results are shown in Figure 4. The peak rate observed during these measurements was 0.08 mr/hr. The M3 production angles during the measurement were 0.46 mr horizontal and 0.9 mr vertical. The measurements were made at the nominal beam height of 48" above the floor of the Detector building (Ref. 6). Both the measurements and the calculation show the muon rates to be highest between the M2 and M3 beam lines. This is due to the deflection of positive muons in the M2 first bend towards the East. Negative muons are produced at a lower rate (due to the lower T production rate) and are deflected in the opposite direction. The rates measured on the west side of the M3 beam are primarily due to negative muons.

Conclusions

"HALO" calculations of muon fluxes due to proton interactions in the Meson center target give results which agree with measurement to within a factor of 3.1. The calculated rates due to $\tau^+ \rightarrow \nu^+$ in Figure 2 are actually higher than the measurements discussed in Appendix 6. This could be due to the fact that the second bend string in the M2 line has not been included and probably spreads the distribution of muons out further (see Appendix 6). Other

inadequacies in the model of the shielding probably have smaller effects.

Muon rates calculated with "HALO" runs at 1000 Gev incident energy are a factor of 6 higher than those for 400 Gev incident, when the Stefanski-White production model is used. The calculated 1000 Gev/400 Gev ratio would probably be an additional factor of 4 higher than this if we had used the production model from the Ferbel 1976 summer study.

If we take our calculation for 1000 GeV primary protons for p^T due to p^T decay and multiply by a safety factor of 2 for the contribution due to kaons, direct production (Ref. 7), and negative muons and a safety factor of 4 for differing x dependences of the production models, we then have a maximum muon flux of 53 x 10**6 per square meter for 10**12 interacting protons—a dosage rate of 11 mr/hr for a 60 second cycle time. This many interacting protons represents 1.8 x 10**12 protons incident on a 16 inch beryllium target (Ref. 8).

The result that we have just quoted neglects the discrepancy between the calculation and the measurement, however. Recognizing this, we could take our calculated rate of 1.4 mr/hr for 1000 GeV incident and 10**12 interacting protons, reduce it by 1/3.1 to account for the discrepancy, boost it by the safety factor of 4 for differing x dependences of the production models, and wind up with a dosage rate of 1.8 mr/hr (8.4 x 10**6 / m^{-2} polse). Our range in the prediction for maximum muon dosage rates is, then, from 1.8 to 11 mr/hr when 1.8 x 10**12, 1000 GeV protons are incident on a 16 inch Be target every 60 sec.

Appendix 1

Attempt to reproduce Ferbel 1976 S.S. results

For pions the production model used was

$$\frac{d\sigma}{dx}(\pi) = 250 e^{-1/x} mb$$

This gives

$$\frac{d\sigma}{d\rho}(\pi^{+}) = \frac{(2.5)(250)}{400} = \frac{-11\times}{1.562} = 1.562$$

The survival fraction for muons with energy above 200 Gev coming from pions produced by 400 Gev protons is

as long as .57 Er < 200. This denominator is correct for the range 0.5 < X $_T$ < 0.8 , which is the range used in the Ferbel summer study.

The fraction of π 's decaying in the first five meters is

$$\left\{1-e^{-\frac{5}{8}c\tau}\right\} \simeq \frac{5}{8c\tau} \left(-\frac{.08944}{E\pi}\right)$$

(where $\text{VCT} \approx$ 56 meters for a 1 GeV T). The relevant cross section is then

$$T_{T} = \int_{0.00}^{320} \left\{ 1.562 \, e^{-11} \, \frac{E_{\Pi}}{400} \right\} \left\{ \frac{E_{\Pi} - 200}{.43E_{\Pi}} \right\} \left\{ \frac{.08944}{E_{\Pi}} \right\}$$

Since Ferbel calculates the effect of re-absorption in the target, we write

$$\int_{0}^{L} e^{-\frac{2}{2}h\rho} \frac{-(L-\frac{2}{2})h\pi}{n dz}$$

$$\int_{0}^{L} e^{-\frac{2}{2}h\rho} \frac{-(L-\frac{1}{2})h\pi}{n dz}$$

$$\int_{0}^{L} e^{-\frac{2}{2}h\rho} \frac{-(L-\frac{1}{2})h\pi}{n dz}$$

where L=20 cm and n=# nuclei cm $^{-3}$ in the target. Continuing

$$\lambda_{1} = \frac{1}{\sqrt{p}n} = \left[38.5(9.01) \cdot \frac{719}{10^{-27}} \frac{6.022}{9.01} \cdot \frac{1.85 \cdot 10^{23}}{9.01} \right]$$

$$\lambda_{p} = \left[26.9 \left(9.01 \right)^{.762} \frac{6.022}{9.01} 1.85 \cdot 10^{23} \right]^{-1}$$

$$\nabla_{EFF} = \frac{-\lambda_p \lambda_{\pi}}{L(\lambda_{\pi} - \lambda_p)} \left[e^{-\frac{t}{\lambda_p}} - e^{-\frac{t}{\lambda_{\pi}}} \right] \nabla_{\tau}'$$

The rate per incident proton is then

$$R = \sigma_{EFF} n L$$

$$= (.66) (1.02 \cdot 10^{-4} \cdot 10^{-27}) (1.236 \cdot 10^{23}) (20)$$

$$= 1.66 \cdot 10^{-7}$$

Multiplying by 3 x 10**12 protons we have the result of 5.0 x 10**5 p+ with energy > 200 GeV per pulse. Since

$$\frac{d\tau}{dx}(\pi^{-}) = \frac{1}{2.5} \frac{d\sigma}{dx}(\pi^{+})$$

we have for the sum of both signs 7.0 x 10**5 p'5 . Ferbel quotes 0.52 x 10**6 muons with energy > 200 GeV, and this is the level to which we reproduce his numbers.

A similar calculation for Kaon parents gave a result of 2.16 x 10**5 muons ($E\rho$ >200 Gev). The corresponding summer study number was 2.3 x 10**5 muons; the agreement is good to within 6%.

Rates for 1000 Gev incident protons were calculated in a similar manner and were a factor of 52 higher than the 400 Gev number for muons with Ey >200 Gev. The corresponding number from the summer study was 60.

Appendix 2

The Stefanski-White Model

The Stefanski-White production model (Ref 2) gives the formula

$$\frac{d^2N}{d\rho d\cos\theta} = ABC^2\rho^2 \left[\exp\left(-8\chi - c\rho \sin\theta\right) \right] \frac{1}{10E_0}$$

where, for T+ production,

$$E_0 = \text{incident energy}$$

$$\chi = PE$$

$$C = C_0 - 1.43 \exp(\chi - \frac{3}{3}\chi^2)$$

$$C_0 = 6$$

$$A = 5$$

$$B = 8$$

If we integrate over an angular range 0 to $\Theta_{\rm M}$, with SiDm 2 $\Theta_{\rm M}$, then

$$\frac{dN}{d\rho} = \frac{A \cdot B}{10 E_0} \left[\exp(-B \times) \right] \left[1 - (1 + C \rho O_m) \exp(-C \rho O_m) \right]$$

For Eo = 400, Tinelastic = 32 mb, and Om = . 01 radians

$$\frac{d\sigma}{d\rho} = .32 \left[\exp(-8 \times) \right] g(\rho)$$

where

$$g(p) = [1 - (1 + .01 cp) exp(-.01cp)]$$

and g(p) ranges from 0.997 to 0.99998985 for p ranging from p=200 to p=360. Therefore

where

The values from this expression are plotted in Fig. 5, along with the 1976 S.S. expression and the fit of Wang (Ref. 9). Figure 6 gives the equivalent results for 1000 GeV incident protons.

Appendix 3

Cross-Check of Calculations

Test no. 1

"HALO" was run with the Stefanski-White production model for and a 5 meter decay path. The number of muons whose energy exceeded 200 GeV was calculated (for 400 GeV protons on target). For 10**12 interacting protons the result was $1.01 \times 10**6$ pt with E > 200 GeV. A direct integration gave $1.0 \times 10**6$ pt and the two results agree to better than the statistical accuracy in the Monte Carlo result.

Test no. 2

"HALO" was run as for test no. 1, except that the incident energy was 1000 GeV. The decay path was again 5 meters. The accepted x range for the π^+ particles was 0.2 to 0.8. The result was 1.76 x 10**7 μ^+ with E μ^+ > 200 GeV. A direct integration gave 1.71 x 10**7 μ^+ and the agreement is acceptable.

Test no. 3

A production model utilizing the Ferbel fit for \mathcal{T}^{\uparrow} production was inserted into "HALO" (Ref 1). When run as in test no. 1 the number of \mathcal{V}^{\uparrow} with E > 200 GeV was 7.5 x 10**5 for 10**12 interacting protons. Direct calculation gave 7.4 x 10**5 \mathcal{V}^{\uparrow} . The level of agreement is acceptable.

Test no. 4

We attempted a comparison of a "HALO" run and the Ferbel S.S. calculation. "HALO" was run with the production formula from test No. 3. The energy ranges were π^+ 200 to 320 GeV, ρ^+ from 200 to 320 GeV. We put in 8.7 meters of steel from 16.4′ to 45′ to represent target load collimators, 38 meters of steel from 145′ to 269′ to represent magnets, and 317 meters of dirt (1041′) from 269′ to 1310′ to represent the berm shielding. The muon patterns were examined at z = 1400′. These absorbers and a fixed dE/dz (no relativistic rise) gave a fixed energy loss of 200 GeV for all muons reaching z = 1400′. The angular range of the π^+ parents was kept to 0 to 0.2 mr, so that the spread at the end was due only to multiple scattering. We multiplied the resulting number of muons by a factor of 60 to account for the missing pion parents at larger angles (This factor came from comparing this run with 0 to 0.2 mr with another that had 0 to 10 mr production angular range).

The "HALO" result is 280 p^+ in a region 1.8 x 1.8 m**2 centered on the o^o beam. The number of interacting protons was 10**12; 2000 $p^+ \rightarrow p^+$ decays were forced to occur in this run. Therefore the fraction of muons falling within a 3 m**2 circle at z=1400' is approximately

$$\frac{280}{2000} \frac{3.0}{3.24} = 13\%$$

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The Ferbel S. S. calculated 60% as the fraction of muons in a central 3 m**2 area. We attempt an explanation of the discrepancy later.

From the "HALO" result the central dosage rate due to p^+ from p^+ would be 0.052 mrem/hr. To this we apply a factor of 2 to account for p^- , p^+ , p^+ parents and a factor of 0.66 to account for absorption in the 20 cm target, and the result is then 0.069 mrem/hr. The equivalent S.S. number is 0.21 mrem/hr (for 10**12 interacting protons). The latter can be multiplied by 0.13/0.60 to account for the difference in the multiple scattering and this would give 0.046 mrem/hr. The level of agreement between the "HALO" result and the corrected S.S. result is a factor of 1.5.

The Ferbel S.S. multiple scattering calculation has a smaller y(r.m.s.) at $z=1400^{\circ}$ than that in the "HALO" calculation because it does not appear to include the effect of the energy loss of the muons. A multiple scattering calculation was done with (Ref. 10) and without energy loss and the difference between the results of the two calculations can account for the difference between the "HALO" calculation and the S.S. calculation regarding the number of muons falling within the central 3 m**2 circle.

Appendix 4

<u>Halo Calculations and Counter Rates in Experiments</u>

"HALO" calculations were first done for the M-center target and the M3 line without the presence of magnetic fields. A model was made for the M3 line for the situation that existed before the "Mesopause" (the Meson area shutdown from September, 1978 to June, 1979) and for the situation that existed after the "Mesopause". A "HALO" run for the pre-"Mesopause" conditions gave fluxes of 1.3 x 10**6 $\slashed{p^+}$ per square meter from $\slashed{m^+}$ decay per 10**12 protons interacting, for the central 1 m x 1 m just outside a beam zone of 16 by 8 square inches. It gave a flux of 1.7 x 10**5 $\slashed{p^+}$ for the 1 m by 1 m area just below the central 1 m by 1 m. These calculated rates were high compared to observed fluxes of charged particles near the M3 beam (Ref. 7) . The 1976 Ferbel Summer Study quotes measured rates near the M3 beam as ranging between 0.6 x 10**5 to 2.7 x 10**5 muons for 3 x 10**12 protons incident on a 8 inch Be target. The conditions are not described, but if they are similar to those of his calculation the area involved would have been within a 4 mr cone about the beam.

The E533 group (Ref. 12) quotes 350000 charged particles per 1.5 x 10**12 protons in an area 50 x 30 square inches, whose lower edge is 8" above the beam. With the beam stop in place this number dropped to 20000. Both are "post-Mesopause" observations and are rates observed with a 16 inch beryllium primary target. The beam stop is made of permanently magnetized material and is a rectangular block 12' long, 4" wide, and 4" high. It is located at z \sim 340 feet. E533 measured rates of 50000 per spill in the same area before the "Mesopause".

Figure 7 shows the p distribution for the "pre-pause" "HALO" calculation. The normalization factor is 15403 muons per unit for 10**12 interacting protons. On this figure the M3 beam axis is at x=0, y=-11.76 inches. In a 50 by 30 square inch area, whose lower edge is 8" above the beam, we find (30)(15403) p on this figure. This converts to 295000 muons for 1.5 x 10**12 protons incident on a 8 inch beryllium target and we see that the number is high compared to the E533 observations. A number more in agreement with their observed rates is gotten by using the output of the "HALO" calculation which includes the effect of 30' of magnet that serves as the M2 first bend and the M3 first magnetic sweeper.

Two further comments are in order. The E533 counters whose rate is being quoted were located about 100 feet downstream of the large aperture 100D40 bending magnet in their experiment. In addition, during the Meson area shutdown, the M3 vacuum pipe from $z=1182^{\prime}$ to 1312' was enlarged from 14" to 36" in diameter and the beam stop was rebuilt to be permanently magnetized.

Appendix 5

More Detail on the "HALO" Runs

The calculation of the muon fluxes for 400 Gev protons striking the M3 target is described here in more detail than in the main text. Limitations imposed by the "HALO" program will be pointed out as we progress.

The production model employed is the Stefanski-White model. from The decay are followed to see if they have sufficient energy to reach the end of the Meson Area and, if so, they are entered into appropriate histograms. The size of the production target is taken to be 0.03 inch by 0.03 inch. Look-up tables are formed for the integrated m+ production formula (see Ref. 1); these are in steps of 10 Gev/c for momenta between 90 and 400 Gev/c and in steps of 0.4 mr for angles between O and 5 milliradian. Next, these tables are modified to account for the fact that we choose to force the decay to occur between z = -5 and 173 feet (see Ref. 1). shows the resulting distribution for the z co-ordinate at which decay of the TT occurs. Entries into the histogram in Fig. 8 are also subjected to the conditions that the pion stayed within apertures before decaying and that the decay muon had sufficient energy to reach z = 1400 feet. Fig. 9 shows a histogram of $dN/d\theta$ for pions at z = 0 for 190 GeV/c. Fig. 10 shows a scatter plot of pat z = 0. Both Figs. 9 and 10 are for only those pions whose decay muons reach the end of the system.

The Meson train-load collimators start 3.3 feet from the target and extend to 50 feet from the target. In the "HALD" input deck these collimators are simulated by 35 collimator elements. This simulation accounts for the tapered apertures — which go from \pm 0.071 by \pm 0.054 inches square to \pm 0.354 by \pm 0.234 inches square (for the M2/M3 beam).

The region $z=45^{\prime}$ to $z=269^{\prime}$ is the Meson Front End hall (see Fig. 1). For this region the "HALO" data deck has aluminum vacuum pipe to 173 $^{\prime}$, has 30 $^{\prime}$ of M2 quadrupole steel represented by steel 13 $^{\prime\prime}$ by 17 $^{\prime\prime}$ (with a 3 $^{\prime\prime}$ dia hole in the center), and has 30 $^{\prime\prime}$ of M2 bend represented by a main ring B2 magnet field map — with a central field of 18.54 kg. (The M2 benders have not been offset in the data deck to represent their positioning so as to allow aperture for both the M2 and M3 beams).

From z=267' to 347' is the FE hall extension. Further downstream are M3 vacuum pipes and enclosures. Berm vacuum pipe 12" in diameter runs between z=415' and 641' and z=695' to 1003'. A pipe 36" in diameter runs between z=1055' and z=1310'. The M3 enclosures are put in as shielding voids, in the shape of simple rectangular boxes. Inside the enclosures we put 6' of steel centered at z=350', 8' of steel centered at z=659', and 9' of steel centered at z=674'. The outside dimensions of this steel ranges

from 12 by 12 square inches at $z=350^{\prime}$ to 36 by 36 square inches at $z=674^{\prime}$. There is a central hole in the steel which ranges from 4 by 2.4 square inches to 6.6 by 4 square inches and is meant to roughly match the M3 beam acceptance.

The dirt berm surrounding the enclosures and berm vacuum pipes has an actual topography which is complex compared to its possible representation by "HALO" shielding configurations. Fig. 11 shows rough contour lines at its top and at its toe. This figure also shows the slope of the trajectory of the M1 and M6 beam lines and it is not surprising that the slope of the toe contour is approximately the same. From this figure we have made rectangular approximations of the dirt berm for use in "HALO". The "tunnel" feature of "HALO" has been used to account for the presence of the dirt berm. From $z=1055^{\circ}$ to $z=1310^{\circ}$ the outside dimensions of this dirt tunnel in the input deck are 168' horizontally and 20' vertically (both centered on the M3 beam). From $z=1310^{\circ}$ to $z=1372^{\circ}$ the horizontal size of the dirt tunnel is increased to 192'.

Changes made to "HALO" by A. Malensk for his studies of muon flux at the 15' bubble chamber were borrowed, in part. These changes modernized the decay constants and radiation lengths (Ref. 12), put in the Stefanski-White production model, included the "density" effect for collision energy losses and pair production in the dE/dz for muons in iron and dirt. Tables 1 and 2 summarize the changes to decay constants and radiation lengths. The dE/dz expressions for iron and soil are give in Table 3. These give values that are reasonable matches to the dE/dz values in Figure 3 of TM-786, "Muon dE/dz and Range Tables for Tevatron Energies: Results for some Shielding Materials", by G. Koizumi.

It is possible to question the simplicity of the shielding model used in these "HALO" studies. This simplicity is imposed by the restrictions inherent in using "HALO". "HALO" is meant to treat a case of one target feeding one beam which travels down one tunnel. We also note here that the magnetic field of the M3 sweeping magnets at z=391' and 675' is not included (these have a polarity opposite to that of the M2 bend at z=247').

The M2 benders centered at $z=374.5^{\circ}$ and $z=391^{\circ}$ are main ring B2-type magnets and total 30° in length. They have outside dimensions 25.25 by 14.25 square inches. The effect of their magnetic field has not been included. The muon distribution pattern at $z=352^{\circ}$ shown in Fig. 13 suggests that the magnetic field of these magnets should have a quite noticeable effect — since many of the muons making it to the end in our study pass through these magnets. The M3 sweeping magnet centered at $z=371^{\circ}$ would have little effect — judging by Fig. 13.

A further improvement to the shielding model would be to account for the M2 tunnel from $z=1003^{\prime}$ to $z=1123^{\prime}$, Judging by Fig. 14 , including the full length of this tunnel would have the effect of removing about 70 $^{\prime}$ of dirt where many muons pass through. Part of

this tunnel is accounted for by the inclusion of the M3 tunnel from z=1003' to z=1055'. The M2 bend magnets in the enclosure from 1003 to 1123 form the third bend point of the M2 beam. Our omission of those magnets should represent a small effect, due to the large vertical size of the muon distribution at that point (see Fig. 14). An improved calculation would account for the second M2 bend point and 70 more feet of enclosure at $z \sim 1100$ feet.

Appendix 6

Small-Scale Survey of Muon Radiation

A small-scale survey of muon background in the Meson Detector Building was made January 19, 1980 in order to have experimentally measured numbers for comparison with those calculated by "HALO" (Ref. 12). The measurements were made between M2 and M3 (4 positions), between M3 and M5 (5 positions), and between M1 and M2 (one position). All were made at the nominal beam height of 48 inches and at a z value = 1386 feet. The portable muon scintillator telescope described as "system 2" in Table 4 was used (Ref. 13).

The instrument was mounted at beam height on a tripod which allowed measurement of relative horizontal and vertical angles. The muon radiation was very directional—as can be seen from figures 15 and 16. At each transverse (x) position the instrument was pointed in the direction of maximum flux and the number of coincidences were recorded for seven beam pulses (the number of protons on target was recorded for the same seven pulses).

The angular widths in Figs. 15 and 16 correspond to the angular cone defined by the two scintillators. This suggests that the muons enter the instrument relatively parallel. Fig. 17 from the "HALO" calculations also shows small angular spreads. We conclude, therefore, that we may take the counts that we observe as the number of muons crossing a 17.8 square centimeter area. By doing so we arrive at the entries in Table 5.

Column 5 of this table is graphed on Fig. 4. The horizontal and vertical production angles of the M3 beam at the time of the measurement were 0.46 mr and \sim 1 mr. The proton beam was going downward and was being bent toward the M1 beam.

REFERENCES

- Ch. Iselin, CERN 74-17, "Halo--A Computer Program to Calculate Muon Halo", 1974. See Fermilab Computer Library Program document PM-33.
- 2. FN-292, R.J. Stefanski and H.B. White, "Neutrino Flux Distributions", May 10, 1976
- 3. The local version originally found on the Cyber 175 had a Sanford-Wang production model built into it. The modification to incorporate the Stefanski-White model is due to A. Malensk. We use here Version 2.2 of "HALO". At CERN the version number has advanced to 2.5.
- 4. The program "LINDA" is partially described in the book, "Particle Accelerator Design Computer Programs", by J. Colonias, Academic Press, 1974, pp. 39-56. The Cyber 175 version was imported from LBL by D. Carey. There is a 370/195 version on the ANL computer, but it was our decision to use the Cyber 175 version in order to "comission" the local "LINDA" capability.
- 5. Section 12.3.3 of the Fermilab Radiation Guide, 3rd edition, P.J. Gollon, Editor, September, 1978 supplies the number 7.8 muons cm⁻² sec⁻¹ = 1 mrem/hr for muons outside a thick shield. This translates to 0.94 x 10**6 muons m⁻² pulse⁻¹ for an accelerator with a repitition rate of a pulse every 12 seconds.
- Since the protons were targetted downward at an angle of 0.9 mr, a height above the floor 15 inches less than the nominal 48 inch beam height would have been a better choice. Figure 2 suggests, however, that the position dependence is not strong enough for 15 inches to make a significant difference.
- 7. The rate for direct production of muons can be crudely estimated relative to pion decay by using a factor of 10**-4 times the pion production rate as the estimate for direct production. The decay fraction for a pion of energy E GeV in L meters is

For a 200 GeV pion the fraction is 10**-4 when L = 1.1 meter.

8. For this calculation we use for the beryllium absorption cross section

$$(38.6)(9.01) = 187.5$$

These numbers are taken from the preprint A.S. Carroll, et. al., Fermilab Pub-78/80-EXP, October, 1978 (Submitted to Physics Letters). The resulting absorption length is 43.1 cm and is larger than the value of 36.7 cm given in the Particle Properties Data Booklet, April, 1978.

- 9. C.L. Wang, Phys. Rev. D 7, 2609 (1973). See also Phys. Rev. D10, 3876 (1974)
- 10. L. Eyges, Phys. Rev. 74, 1534 (1948). This paper was brought to my attention by Fermilab TM-261, "Muon Shielding: Multiple Coulomb Scattering of Muons with Energy Loss", by D. Theriot. The reader of either of these papers is advised to check the equations for an error of a factor of two.
- 11. The measument equipment is, of course, insensitive to the sign of the muons. A great deal of sign separation occurs at the first bending magnets. The positive muons are mostly bent to the east of the M3 line. They are what has been calculated and they are responsible for the highest points seen in Figure 4.
- 12. private communication, Bruce Winstein.
- 13. The muon measurements were performed in conjunction with W. Baker.
- 14. The portable scintillator telescope is described in section 4.10 of the Fermilab Radiation Guide, op. cit. B.J. Holt was very helpful in arranging the loan of this instrument.

A. Malensk

Particle

	Program (Version 2.	change 2)	Data Booklet April, 1978
	c 1	~ values (meters)	
T Decay	7. 64	7.8	7. 804
T Decay	3. 67	3. 71	3. 709
		Branching Ratio	
K-> Y2	0. 58	0. 63	0. 635

HALO

Table 1
Change in Decay Constants

Radiation Lengths (meters)

Material	HALO Program (Version 2.2)	A. Malensk change	Particle Data Booklet April, 1978
AIR	312. 4	312. 4	300. 5
FE	. 018	. 0177	. 0176
cu	0147	0147	. 0143
DIRT	. 15	165	
Be	. 338	. 338	. 353
РЬ	. 0051	. 0051	. 0056
Al	. 0886	. 0886	. 089

Table 2
Change in Radiation Lengths

p range

dE/dz for Iron (Gev/meter)

< 1

1.16

100 1.26+0.12 ln (p)

100 ≤ p < 400

7.72-2.69 ln (p) + 0.31 (ln(p))

400

32.08-10.61 ln (p) + 0.954 $\left(\ln(p)\right)^2$

dE/dz for Iron

p range

dE/dz for Soil (Gev/meter)

< 1

. 46

1 < p < 100

 $0.35 + 0.0287 \ln (p)$

100 < p < 400

1.51-0.434 ln (p) + 0.0466 $(\ln(p))^2$

400 < p < 1000

3. 3854-1. 184 ln (p) + 0. 107 $\left(\ln(p)\right)^2$

dE/dz for Soil (Gev/meter)

These

detectors and the asociated electronics can be carried by one person. Their LED scalers can be set to accumulate counts in one of four modes: coincidences between the two detectors, singles in either one of the two detectors, or random (i.e., delayed) coincidences between the front and back detectors.

Scintillator Diameter	4.8 cm
Scintillator Area	17.8 cm ²
Scintillator Spacing	53 cm
Half-Angle of Cone of Sensitivity	45 mrad (2.6°)

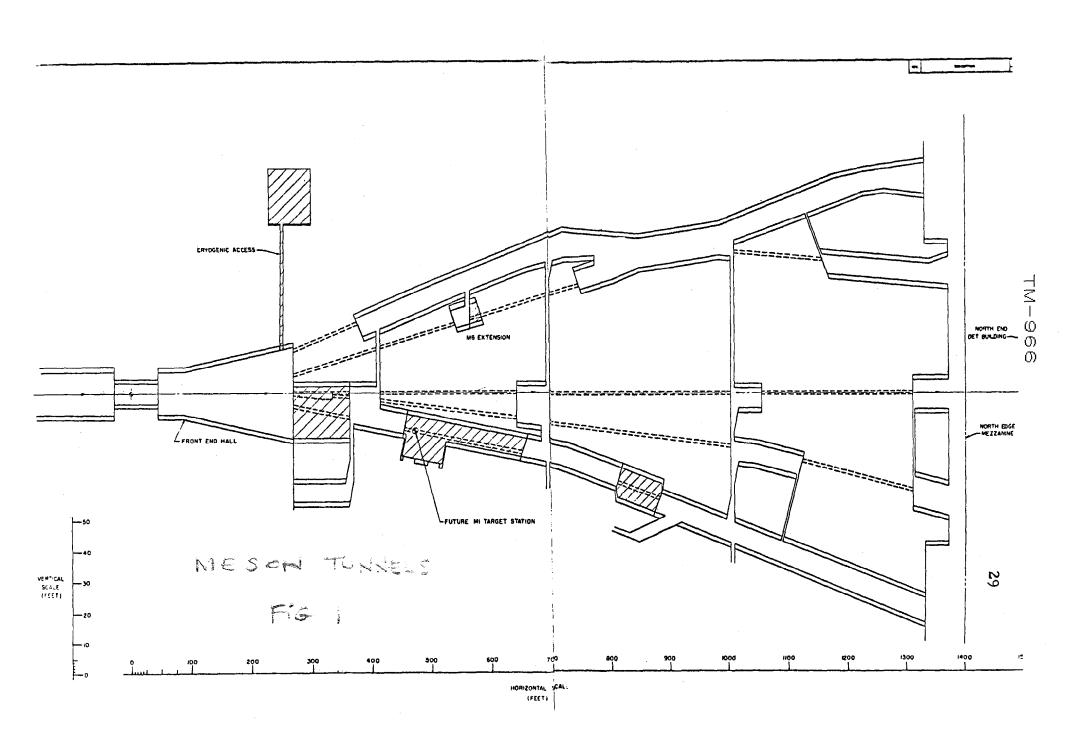
Properties of Radiation Physics "Muon Finder"

Table 4

x value (inches)	No. of counts		is muons	cting
-628	75	3. 7	1. 7	0.018
-296	2423	25. 9	7. 8	0.084
-234	972	27	3. 0	0. 032
-177	423	26	1.4	0.015
-124	1014	25	3. 4	0. 036
123	301	22. 3	1.1	0.012
196	254	25. 4	0.8	0.009
264	423	24. 5	1.4	0.015
325	332	25. 7	1. 1	0.012
411	95	25	О. З	0. 0034

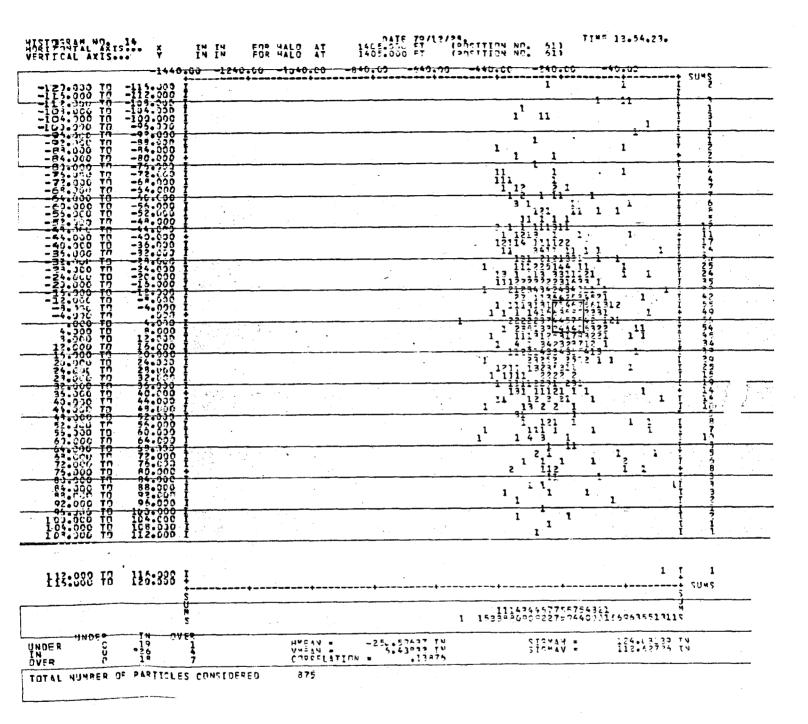
Table 5

Muon Flux Measurements at Nominal Beam Height

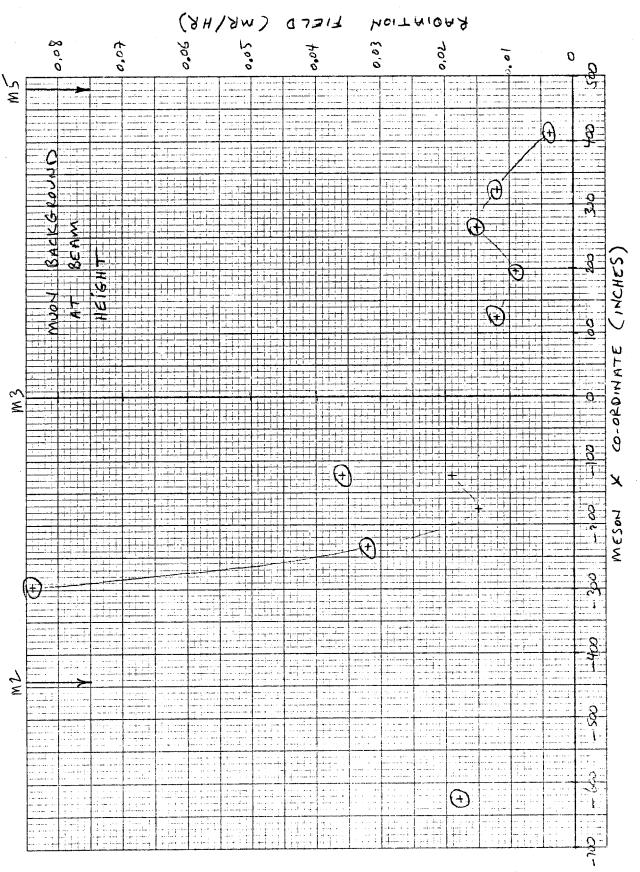


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-112.000	TO TO	-108.000 -104.000	I	1 1		11	3	
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-32.000 -28.030	ŤŌ	-24.000	1	133 311 121 2245 21 2441 1 11121 2 5112 2 2 3 2 4 4 4	3231 11	2	I 31	
-24.000 -20.000	TO TO	-20.000 -16.000	1	4321	[5 3241	1	I 25	
-16.000 -12.000	TO	-12.000 -8.000	1	1,11121	52431	11	1 29 1 25 1 31	
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24.000 28.000	TO	28.000 32.000	1_1_1	12 3 32 4 22154 1 1 32 114112	13211	2 1	I 27	
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68.000	TO ·	72.000	<u> I</u> 1	21		1 .	Ī 11	
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80.000 84.000	TO	84.000 88.000		1 1 1	2 1	1 1	I 6	
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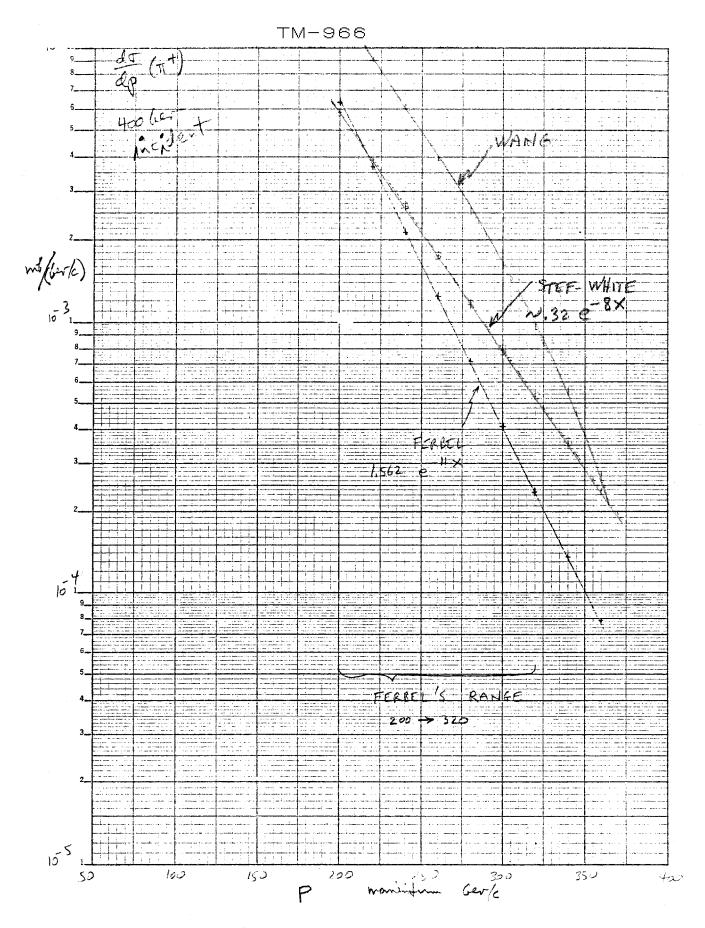
MUON DISTRIBUTION 400 GEV INCIDENT FIGURE 2



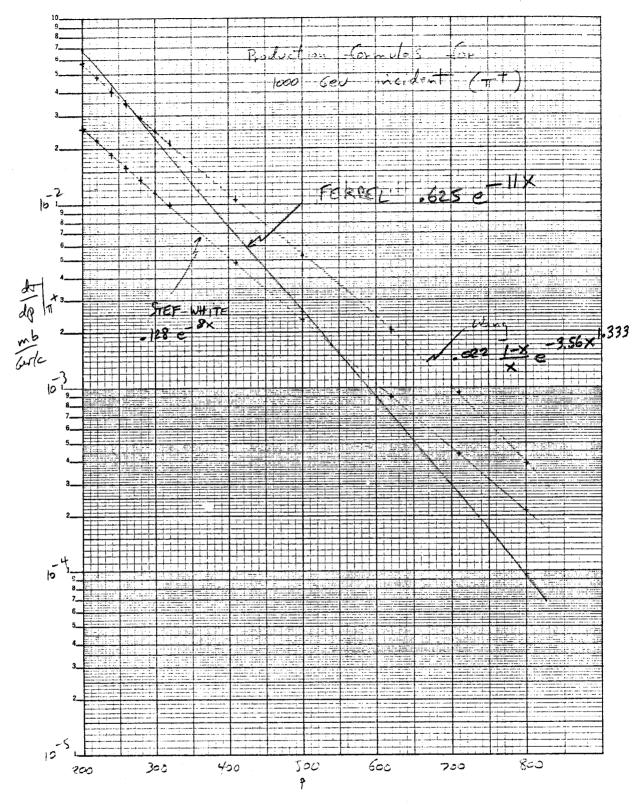
MUON DISTRIBUTION
1000 GEV INCIDENT
FIG. 3



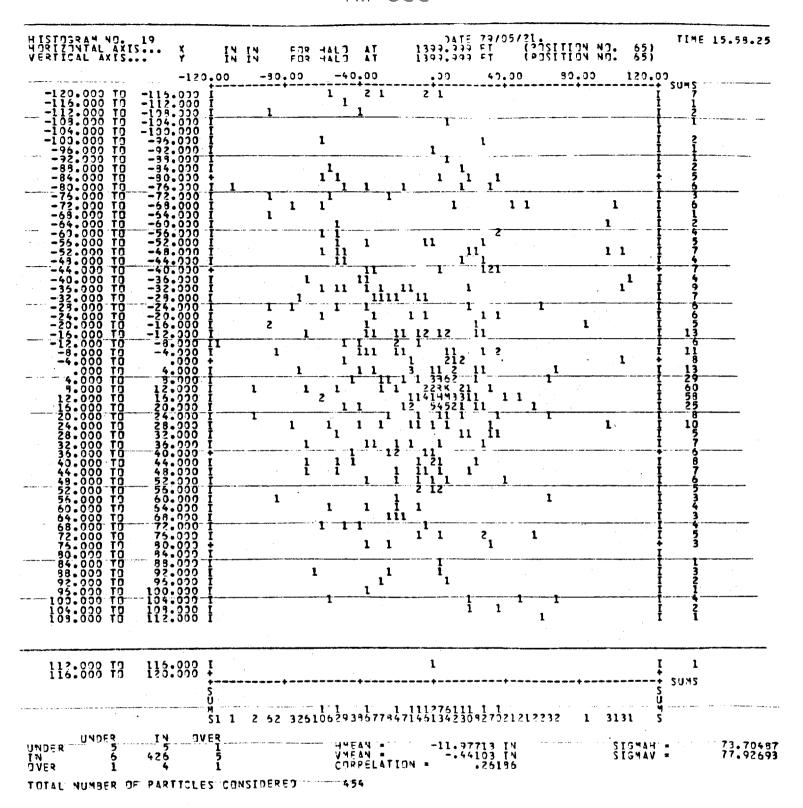
MUDA MEASUREMENTS FIGURE 4



PRODUCTION CURVE too GEV FIGURE 5



PRODUCTION CURVES
1000 GEV
FIGURE 6



PRE- PAUSE 400 GEV INCIDENT FIGURE 7

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ESS THAN	0.000	0							
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40.000 TO 45.000 TO 50.000 TO	45.000 50.000 55.000 60.000	17 21	XX7 X7 XX1 XX1						
60.000 TO 65.000 TO 70.000 TO 75.000 TO	65.CCO 70.000 75.000 80.000	21 25 26 30 27	7XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX						
80.000 TO 85.000 TO 90.000 TO 95.000 TO	85.000 90.000 95.000 100.000	30 27 20 25 21 15 19	XX5 XX1 X5				· · ·		•
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DECAY POINT DISTRIBUTION

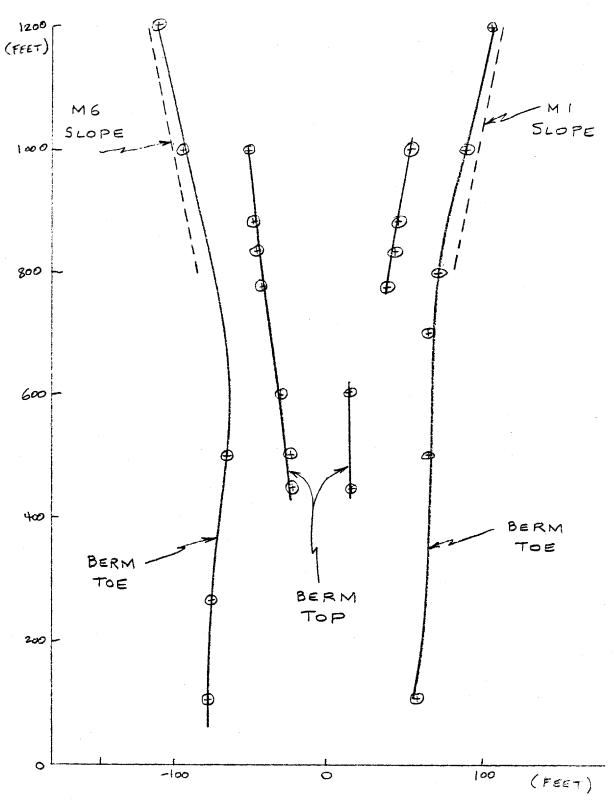
Fig. 8

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2.400 TD 2.600 TD 2.800 TD 3.000 TD	2.600 2.800 3.000 3.200	0 1 X 0 0	
3.200 TD 3.400 TD 3.600 TD 3.800 TD	3.400 3.600 3.800 4.000	1 X 1 X	
4.000 TO 4.200 TO 4.400 TO 4.600 TO	4.200 4.400 4.600 4.800	2 XX 1 X 0	
4.800 TD 5.000 TD 5.200 TD 5.400 TD	5.000 5.200 5.400 5.600	0 0 0	
5.600 TD 5.800 TD 6.000 TD 6.200 TD	5.800 6.000 6.200 6.400	0 0 0 0	
6.400 TO 6.600 TO 6.800 TO 7.000 TO	6.600 6.800 7.000 7.200	0 0 0	
7.200 TO 7.400 TO 7.600 TO 7.800 TO	7.400 7.600 7.800	0	
8.000 TD 8.200 TD 8.400 TD 8.600 TD	8.20C 8.400 8.600 8.800	0 0 0	
8.800 TD 9.000 TD 9.200 TD 9.400 TD	9.000 9.200 9.400 9.600	0	
9.600 TO 9.800 TO	9.800 10.000	0	
GREATER THAN	10.000	0	
		.S. HALF WIDTH =	.91281 MR
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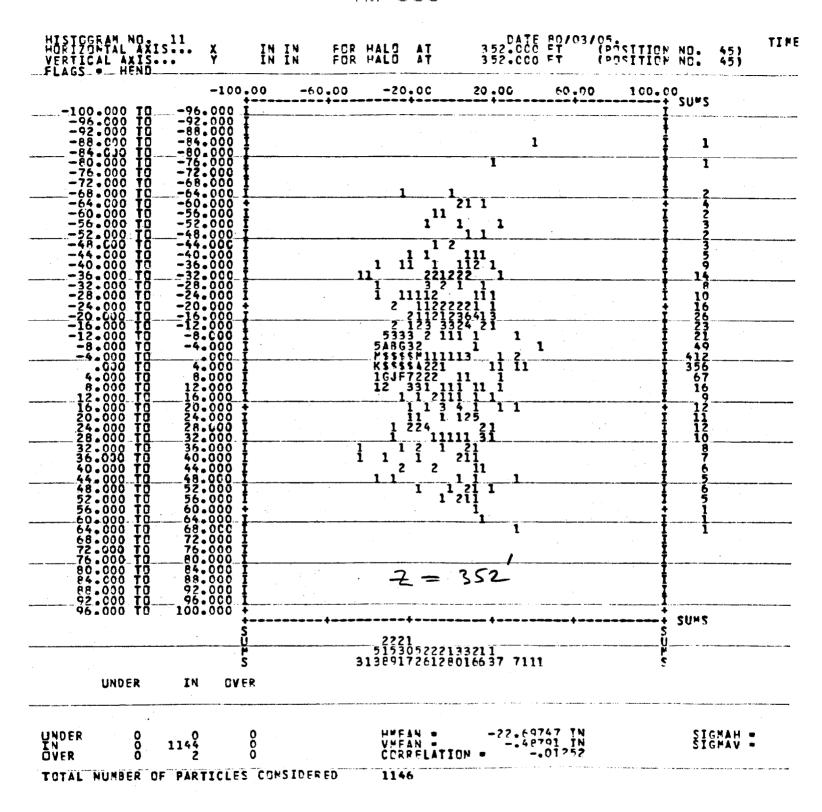
FIGURE 9

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BERM. CONTOURS

Fig. 11



MUON DISTRIBUTION

400 GEU INCIDENT

FIGURE 12

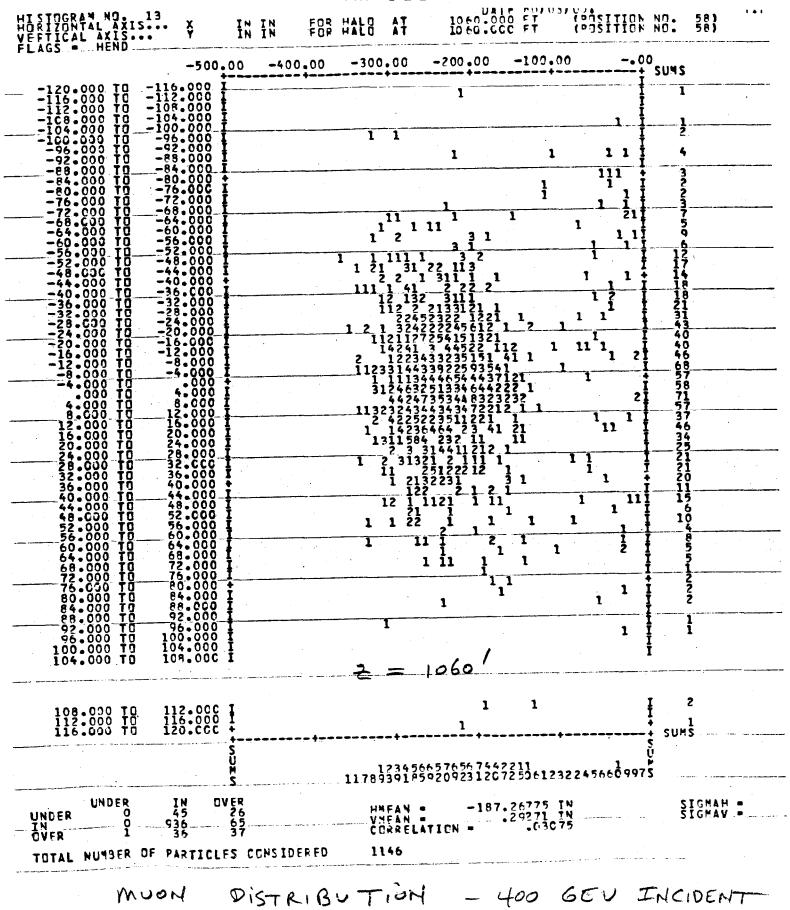
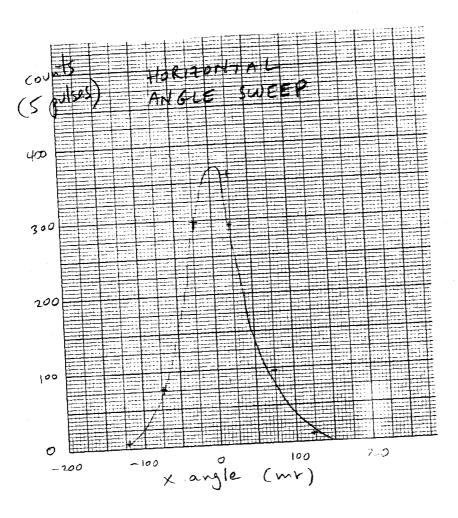
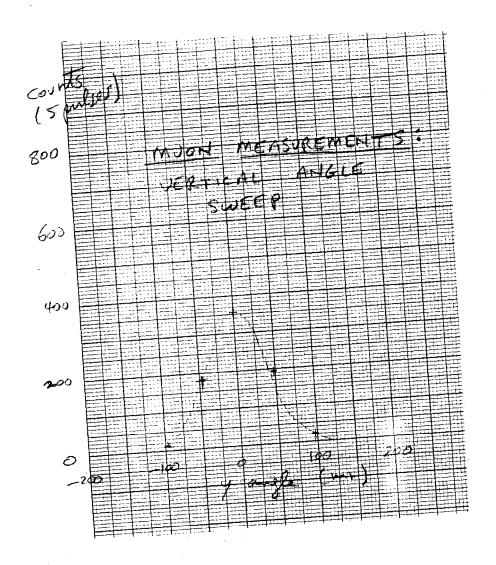


Fig. 13

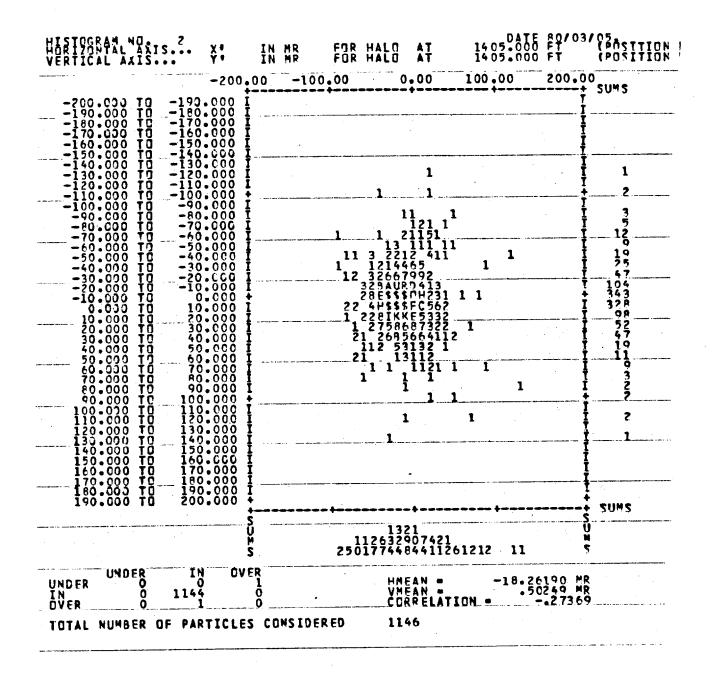


ANGULAR WIDTH MUON MEASUREMENT FIG. 14



ANGULAR WIDTH
MUON MEASURENENTS

Fig. 15



X', Y' SCATTER PLOT MUON CALCULATION 400 GEU INCIDENT FIGURE 16